Laplacian Sum-Eccentricity Energy of a Graph

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Abstract

We introduce the Laplacian sum-eccentricity matrix $\mathbf{LS}_{\mathbf{e}}$ of a graph G, and its Laplacian sum-eccentricity energy $LS_eE = \sum_{i=1}^{n} |\eta_i|$, where $\eta_i = \zeta_i - \frac{2m}{n}$ and where $\zeta_1, \zeta_2, \ldots, \zeta_n$ are the eigenvalues of $\mathbf{LS}_{\mathbf{e}}$. Upper bounds for LS_eE are obtained. A graph is said to be twinenergetic if $\sum_{i=1}^{n} |\eta_i| = \sum_{i=1}^{n} |\zeta_i|$. Conditions for the existence of such graphs are established.

Keywords: Sum-eccentricity eigenvalues, sum-eccentricity energy, Laplacian sum-eccentricity matrix, Laplacian sum-eccentricity energy.

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1. Introduction

Let G be a simple connected graph with vertex set $\mathbf{V}(G)$ and edge set $\mathbf{E}(G)$, of order $|\mathbf{V}(G)| = n$ and size $|\mathbf{E}(G)| = m$. Let $\mathbf{A} = (a_{ij})$ be the adjacency matrix of G. The eigenvalues $\lambda_1, \lambda_2, \ldots, \lambda_n$ of \mathbf{A} are the eigenvalues of the graph G [6]. Since \mathbf{A} is a symmetric matrix with zero trace, these eigenvalues are real with sum equal to zero. The energy of the graph G is defined as the sum of the absolute values of its eigenvalues [10, 16]:

$$E(G) = \sum_{i=1}^{n} |\lambda_i|.$$

After the introduction of the graph–energy concept in the 1970s [10], several other "graph energies" have been put forward and their mathematical properties extensively studied; for details see the recent monograph [12] and the survey [11].

In the last few years, a whole class of graph energies was conceived, based on the eigenvalues of matrices associated with a particular topological index. Thus, let TI be a topological index is of the form

$$TI = TI(G) = \sum_{v_i v_j \in \mathbf{E}(G)} F(v_i, v_j)$$

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where F is a pertinently chosen function with the property F(x, y) = F(y, x). Then a matrix **TI** can be associated to TI, defined as

$$(\mathbf{TI})_{ij} = \begin{cases} F(v_i, v_j) & \text{if } v_i v_j \in \mathbf{E}(G) \\ 0 & \text{otherwise.} \end{cases}$$

If $\tau_1, \tau_2, \ldots, \tau_n$ are the eigenvalues of the matrix **TI**, then an "*energy*" can be defined as

$$E_{TI} = E_{TI}(G) = \sum_{i=1}^{n} |\tau_i|.$$
 (1)

The most extensively studied such graph energy is the *Randić energy* [2,3,7,12], based on the eigenvalues of the Randić matrix **R**, where

$$(\mathbf{R})_{ij} = \begin{cases} \frac{1}{\sqrt{d_i \, d_j}} & \text{if } v_i v_j \in \mathbf{E}(G) \\ 0 & \text{otherwise} \end{cases}$$

and where d_i is the degree of the *i*-th vertex of *G*. In an analogous manner the harmonic energy [14], *ABC* energy [9], geometric–arithmetic energy [23], Zagreb energy [15], and sum-eccentricity energy [22,26] were put forward.

To any energy E_{TI} of the form (1), a "Laplacian energy" LE_{TI} can be associated, defined as

$$LE_{TI} = LE_{TI}(G) = \sum_{i=1}^{n} \left| \theta_i - \frac{2m}{n} \right|$$
(2)

where $\theta_1, \theta_2, \ldots, \theta_n$ are the eigenvalues of the matrix $\mathbf{LTI} = \mathbf{D} - \mathbf{TI}$, and where $\mathbf{D} = \mathbf{D}(G)$ is the diagonal matrix of vertex degrees.

The first such Laplacian energy, based on the adjacency matrix \mathbf{A} , was introduced in 2006 [13] and its theory is nowadays elaborated in full detail, see [12]. It is worth noting that this Laplacian energy found interesting engineering applications in image processing [18, 25, 27]. Bearing this in mind, it is purposeful to study other Laplacian graph energies. Some recent studies along these lines are [1, 5, 8, 20, 21].

In this paper we study the Laplacian version of the sum-eccentricity energy. In order to define it, we need some preparations.

The distance d(u, v) between two vertices u and v in a (connected) graph G is the length of a shortest path connecting u and v [4]. The eccentricity of a vertex $v \in \mathbf{V}(G)$ is $e(v) = max\{d(u, v) : u \in \mathbf{V}(G)\}$. The radius of G is $r(G) = min\{e(v) : v \in \mathbf{V}(G)\}$, whereas the diameter of G is $d(G) = max\{e(v) : v \in \mathbf{V}(G)\}$. Hence $r(G) \leq e(v) \leq d(G)$, for every $v \in \mathbf{V}(G)$.

In this paper, we denote by K_n , $K_{a,b}$, $K_{1,a}$, C_n , and P_n the complete graph, complete bipartite graph, star, cycle, and path, respectively.

The sum-eccentricity matrix of a graph G is denoted by $\mathbf{S}_{\mathbf{e}}(G)$ and defined as $\mathbf{S}_{\mathbf{e}}(G) = (s_{ij})$ [22, 26], where

$$s_{ij} = \begin{cases} e(v_i) + e(v_j) & \text{if } v_i v_j \in E \\ \\ 0 & \text{otherwise.} \end{cases}$$

If $\mu_1, \mu_2, \ldots, \mu_n$, are the eigenvalues of $\mathbf{S}_{\mathbf{e}}(G)$, then the sum-eccentricity energy is

$$ES_e(G) = \sum_{i=1}^n |\mu_i|.$$

Definition 1.1. Let G be a graph of order n and size m. The Laplacian sumeccentricity matrix of G, denoted by $\mathbf{LS}_{\mathbf{e}}(G) = (\ell_{ij})$, is defined as

$$\mathbf{LS}_{\mathbf{e}}(G) = \mathbf{D}(G) - \mathbf{S}_{\mathbf{e}}(G) \,.$$

The Laplacian sum-eccentricity spectrum of G, consisting of $\zeta_1, \zeta_2, \ldots, \zeta_n$, is the spectrum of the Laplacian sum-eccentricity matrix. This leads us to define the Laplacian sum-eccentricity energy of a graph G as

$$LS_e(G) = \sum_{i=1}^n \left| \zeta_i - \frac{2m}{n} \right|.$$
(3)

If, in addition, we define the auxiliary quantity η_i as

$$\eta_i = \zeta_i - \frac{2m}{n}$$

then

$$LS_e E(G) = \sum_{i=1}^n |\eta_i|.$$

Lemma 1.2. Let G be an (n,m)-graph. Then

$$\sum_{i=1}^n \zeta_i = 2m \,.$$

Proof.

$$\sum_{i=1}^{n} \zeta_i = trace(\mathbf{LS}_{\mathbf{e}}(G)) = \sum_{i=1}^{n} \ell_{ii} = \sum_{i=1}^{n} d_i = 2m.$$

Theorem 1.3. The Laplacian sum-eccentricity energy of the complete graph K_n is

$$LS_e E(K_n) = 4(n-1).$$

Proof. Recalling that the eccentricity of any vertex of K_n is unity, directly from the definition of the Laplacian sum-eccentricity matrix, we calculate that

$$Spec(\mathbf{LS}_{\mathbf{e}}(K_n)) = \begin{bmatrix} -d & d+2\\ 1 & n-1 \end{bmatrix}$$

where d = n - 1 is the degree of any vertex of K_n . Using the fact that $\frac{2m}{n} = n - 1 = d$, we get by Eq. (3)

$$LS_e E(K_n) = |-d-d| + |d+2-d| + \dots + |d+2-d|$$

= $2d + 2(n-1) = 4(n-1).$

2. Bounds for Laplacian Sum-Eccentricity Energy

Theorem 2.1. Let G be an (n,m)-graph. Then

$$LS_e E(G) \le \sqrt{n\left(\sum_{i=1}^n \sum_{j=1}^n \ell_{ij}^2 - \frac{4m^2}{n}\right)}.$$
 (4)

Proof. We have

$$\sum_{i=1}^{n} \eta_i^2 = \sum_{i=1}^{n} \left(\zeta_i - \frac{2m}{n}\right)^2 = \sum_{i=1}^{n} \left(\zeta_i^2 - \frac{4m}{n}\zeta_i + \frac{4m^2}{n^2}\right)$$
$$= \sum_{i=1}^{n} \zeta_i^2 - \frac{4m}{n} \sum_{i=1}^{n} \zeta_i + \frac{4m^2}{n}.$$

By Lemma 1.2, $\,$

$$\sum_{i=1}^{n} \eta_i^2 = \sum_{i=1}^{n} \zeta_i^2 - \frac{8m^2}{n} + \frac{4m^2}{n} = \sum_{i=1}^{n} \zeta_i^2 - \frac{4m^2}{n}.$$

Using the Cauchy–Schwarz inequality

$$\sum_{i=1}^n |\eta_i| \le \sqrt{n \sum_{i=1}^n \eta_i^2}$$

we get

$$\sum_{i=1}^{n} |\eta_i| \le \sqrt{n\left(\sum_{i=1}^{n} \zeta_i^2 - \frac{4m^2}{n}\right)}.$$

On the other hand,

$$\sum_{i=1}^n \zeta_i^2 = trace(\mathbf{LS^2_e}(G)) = \sum_{i=1}^n \sum_{j=1}^n \ell_{ij}^2$$

and inequality (4) follows.

It should be noted that inequality (4) is just a variant of the classical McClelland's upper bound for ordinary graph energy [16, 17].

Corollary 2.2. Let G be an r-regular graph. Then

$$\sum_{i=1}^{n} \eta_i^2 = \sum_{i=1}^{n} \sum_{j=1}^{n} \ell_{ij}^2 - nr^2.$$

Example 2.3. If $G \cong K_n$, then

$$\sum_{i=1}^{n} \eta_i^2 = 4n(n-1).$$

If $G \cong K_{a,b}$. Then

$$\sum_{i=1}^{n} \eta_i^2 = ab \left(a + b + 32 - \frac{4ab}{a+b} \right) \,.$$

In particular, for $G \cong K_{1,a}$, with n = a + 1:

$$\sum_{i=1}^{n} \eta_i^2 = a^2 \left(1 - \frac{4}{a+1} \right) + 19a \,.$$

In what follows we derive another upper bound for the Laplacian sum-eccentricity energy using Weyl's inequality for matrices.

Theorem 2.4. (Weyl's inequality) [19] Let \mathbf{X} and \mathbf{Y} be Hermitian $n \times n$ matrices. If for $1 \leq i \leq n$, $\lambda_i(\mathbf{X})$, $\lambda_i(\mathbf{Y})$, $\lambda_i(\mathbf{X}+\mathbf{Y})$ are the eigenvalues of \mathbf{X} , \mathbf{Y} , and $\mathbf{X}+\mathbf{Y}$, respectively, then

$$\lambda_i(\mathbf{X}) + \lambda_n(\mathbf{Y}) \le \lambda_i(\mathbf{X} + \mathbf{Y}) \le \lambda_i(\mathbf{X}) + \lambda_1(\mathbf{Y})$$

The matrices $\mathbf{LS}_{\mathbf{e}}(G)$, $\mathbf{S}_{\mathbf{e}}(G)$, and $\mathbf{D}(G)$ are all Hermitian $n \times n$ matrices. In addition, we use the facts that the eigenvalues of the diagonal matrix are the entries in the diagonal, and that the energy of a matrix \mathbf{X} is equal to the energy of $-\mathbf{X}$. We thus arrive at:

Theorem 2.5. Let G be an (n,m)-graph with maximal vertex degree Δ . Then

$$LS_e E(G) \le ES_e(G) + k\,\Delta + \frac{2m}{n}(n-k) \tag{5}$$

where $k = |\{\zeta_i : \zeta_i \ge 2m/n\}|.$

Proof. Let $\zeta_1 \geq \zeta_2 \geq \cdots \geq \zeta_n$ be the Laplacian sum-eccentricity eigenvalues, $\mu_1 \geq \mu_2 \geq \cdots \geq \mu_n$ be the sum-eccentricity eigenvalues and $\rho_1 \geq \rho_2 \geq \cdots \geq \rho_n$ be the eigenvalues of the degree matrix. We assume that $1 \leq k \leq r \leq n$. Using Theorem 2.4 we get

$$\mu_i + \rho_n \le \zeta_i \le \mu_i + \rho_1$$

Since $\rho_n \ge 0$,

$$\mu_i \le \zeta_i \le \mu_i + \rho_1 \,.$$

Since $2m/n \ge 0$,

$$\mu_i - \frac{2m}{n} \le \zeta_i - \frac{2m}{n} \le \mu_i + \rho_1.$$

Now we have to distinguish between two cases.

Case 1: If $\zeta_i - \frac{2m}{n} \ge 0$, then

$$\left|\zeta_i - \frac{2m}{n}\right| \le \mu_i + \rho_1 \,.$$

If there are $k~\zeta_i{'}{\rm s},$ satisfy this condition, then

$$\sum_{i=1}^{k} \left| \zeta_i - \frac{2m}{n} \right| \le \sum_{i=1}^{k} (\mu_i + \rho_1) \le \sum_{i=1}^{k} |\mu_i| + k\rho_1.$$
 (6)

Case 2: If $\zeta_i - \frac{2m}{n} \leq 0$, then

$$\left|\zeta_i - \frac{2m}{n}\right| \le \left|\mu_i - \frac{2m}{n}\right|.$$

If we have, $\mu_i \leq 0$ for i = k + 1, ..., r and $\mu_i \geq 0$ for i = r + 1, ..., n. Then

$$\sum_{i=k+1}^{n} \left| \zeta_{i} - \frac{2m}{n} \right| \leq \sum_{i=k+1}^{r} \left| \mu_{i} - \frac{2m}{n} \right| + \sum_{i=r+1}^{n} \left| \mu_{i} - \frac{2m}{n} \right|$$
$$= \sum_{i=k+1}^{r} \frac{2m}{n} + \sum_{i=k+1}^{r} |\mu_{i}| + \sum_{i=r+1}^{n} \frac{2m}{n} - \sum_{i=r+1}^{n} |\mu_{i}|.$$
(7)

Combining the relations (6) and (7), we get

$$\sum_{i=1}^{n} \left| \zeta_{i} - \frac{2m}{n} \right| \leq \sum_{i=1}^{k} |\mu_{i}| + k\rho_{1} + \sum_{i=k+1}^{r} |\mu_{i}| - \sum_{i=r+1}^{n} |\mu_{i}| + \frac{2m}{n} (n-k)$$
$$\leq ES_{e}(G) + k\rho_{1} + \frac{2m}{n} (n-k)$$

from which (5) follows straightforwardly.

Corollary 2.6. If the graph G is r-regular, then

$$LS_e E(G) \le ES_e(G) + nr.$$
(8)

Proof. From Theorem 2.5, we have

$$LS_e E(G) \le ES_e(G) + kr + \frac{2m}{n}(n-k).$$

Since, in addition, for an r-regular graph, 2m/n = r,

$$LS_e E(G) \le ES_e(G) + r(k+n-k)$$

and inequality (8) follows.

Lemma 2.7. [24] For the complete bipartite graph $K_{a,b}$, the sum-eccentricity energy is $ES_e(K_{a,b}) = 8\sqrt{ab}$.

Corollary 2.8. For the complete bipartite graph $K_{a,b}$,

$$LS_e E(K_{a,b}) \le 8\sqrt{ab} + k \max\{a,b\} + \frac{2ab}{a+b}(a+b-k).$$
 (9)

Proof. For $K_{a,b}$, 2m/n = 2ab/(a+b). Using Lemma 2.7, we get (9) from (5).

3. Twinenergetic Graphs

In this section, we point out a remarkable feature of Laplacian sum-eccentricity energy.

Definition 3.1. Let G be a graph of order n, and let ζ_i , i = 1, 2, ..., n, be its Laplacian sum-eccentricity eigenvalues. We say that G is twinenergetic if

$$LS_e E(G) = \sum_{i=1}^n |\zeta_i|.$$

The above definition means that

$$\sum_{i=1}^{n} \left| \zeta_i - \frac{2m}{n} \right| = \sum_{i=1}^{n} \left| \zeta_i \right|.$$
 (10)

The number of positive eigenvalues and negative eigenvalues (including their multiplicities) are denoted by $\zeta^+(G)$ and $\zeta^-(G)$, respectively. For the sake of simplicity, we assume that there are no zero Laplacian sum-eccentricity eigenvalues, i.e., that $\zeta^+(G) + \zeta^-(G) = n$.

Theorem 3.2. A graph G is a Laplacian sum-eccentricity twinenergetic if it satisfies the following two conditions:

$$\zeta_i(G) \geq \frac{2m}{n}, \ i = 1, 2, \dots, \zeta_i^+(G).$$
 (11)

$$\zeta^+(G) = \zeta^-(G). \tag{12}$$

Proof. Let $\zeta^+(G) = r$, where $1 \le r \le n$. Then

$$\sum_{i=1}^{n} \left| \zeta_{i} - \frac{2m}{n} \right| = \sum_{i=1}^{r} \left| \zeta_{i} - \frac{2m}{n} \right| + \sum_{i=r+1}^{n} \left| \zeta_{i} - \frac{2m}{n} \right|$$
$$= \sum_{i=1}^{r} \left| \zeta_{i} - \frac{2m}{n} \right| + \sum_{i=r+1}^{n} \left| \zeta_{i} \right| + \frac{2m}{n} (n-r) \,. \tag{13}$$

Let k be the number of eigenvalues satisfying the condition $\zeta_i(G) \geq \frac{2m}{n}$. Then, $1 \leq k \leq r$, and

$$\sum_{i=1}^{r} \left| \zeta_{i} - \frac{2m}{n} \right| = \sum_{i=1}^{k} \left(\zeta_{i} - \frac{2m}{n} \right) + \sum_{i=k+1}^{r} \left(\frac{2m}{n} - \zeta_{i} \right)$$
$$= \sum_{i=1}^{k} |\zeta_{i}| - \frac{2m}{n} k - \sum_{i=k+1}^{r} |\zeta_{i}| + \frac{2m}{n} (r - k). \quad (14)$$

Substituting (14) back into (13) yields

$$\sum_{i=1}^{n} \left| \zeta_i - \frac{2m}{n} \right| = \sum_{i=1}^{k} |\zeta_i| + \sum_{i=r+1}^{n} |\zeta_i| - \sum_{i=k+1}^{r} |\zeta_i| + \frac{2m}{n} (n-2k)$$
$$= \sum_{i=1}^{n} |\zeta_i| - 2\sum_{i=k+1}^{r} |\zeta_i| + \frac{2m}{n} (n-2k).$$

If the condition (11) is obeyed, i.e., if k = r, then

$$\sum_{i=k+1}^r |\zeta_i| = 0.$$

If, in addition, also the condition (12) is obeyed, i.e., 2r = n, then

$$\frac{2m}{n}(n-2k) = 0\,.$$

Thus, if both conditions (11) and (12) are satisfied, then the relation (10) holds, i.e., the graph G is twinenergetic. $\hfill \Box$

It now remains to see if twinenergetic graphs exist at all. That such graphs do exist is verified by the following examples.

Example 3.3. The paths P_2 , P_4 , P_8 , and P_{16} are twinenergetic graphs.

For P_2 , direct calculations gives $\zeta_1 = 3$, $\zeta_2 = -1$, and 2m/n = 1. Therefore,

$$\sum_{i=1}^{2} |\zeta_i| = 4 \quad \text{and} \quad \sum_{i=1}^{2} \left| \zeta_i - \frac{2m}{n} \right| = 4.$$

For P_4 we get $\zeta_1 = 9.0902$, $\zeta_2 = 4.7202$, $\zeta_3 = -2.0902$, $\zeta_4 = -5.7202$, and 2m/n = 6/4. Therefore,

$$\sum_{i=1}^{4} |\zeta_i| = \sum_{i=1}^{4} \left| \zeta_i - \frac{2m}{n} \right| = 21.6208.$$

The cases P_8 and P_{16} are verified analogously.

Finding more examples of twinenergetic graphs, as well as their complete structural characterization remains a task for the future.

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